2D FE and 2DOF Simulations of Ground Shock Experiments – Reflection Pressure Time History Dependency due to the Charge's and Structure's Properties

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This paper analyses, by using 2D axial symmetry Finite Element (FE) with Autodyn, the structural response of a well-defined structure; a suspended piston-spring system buried in sand subjected to ground shock from an explosive charge. The parameters varied in the simulations were charge size, charge distance, reflection area of the piston, piston mass, and spring stiffness. Earlier experiments from 1980s, conducted by S. Hultgren, FortF, showed that the reflection pressure over time was dependent on the mass and stiffness of the structure. Here, some more parameters were varied in simulations to see how well a two-degree-of-freedom model (2DOF) can capture the main behaviour of the structural response. The first aim of the FE simulations was to better understand the physics of the observed experimental results, which has been confirmed in earlier studies [1], [2]. Based on this, the second aim was to find a methodology that can use simplified relationships for ground shock prediction, from e.g. ConWep, in combination with simplified models such as 2DOF, to predict the structural response of e.g. a buried concrete wall.

The FE simulation models were generated in Autodyn-2D, where the sand was modelled with Euler cells and the piston, spring, and cylinder were modelled with Lagrange elements. The sand was modelled with an Equation of State (EOS) designed for porous soils. The simple 2DOF model confirms some of the main behaviour found in the FE results and experiments, such as the initial collision effect only depends on mass and that total spring deflection gave reasonable agreement with simulations.

INTRODUCTION

The Swedish Civil Contingencies Agency (MSB) is responsible for the building regulations of the Swedish civil defence shelters. There are specific regulations for how the defence shelters are planned, built, equipped, and maintained [3]. One of many regulations state what load level the shelters should be able to withstand: "The effect of a pressure wave corresponding to that produced by a 250 kg GP-bomb with 50 weight percent TNT which burst freely outside at a distance of 5.0 meters from the outside of the shelter during free pressure release". However, many of the shelters are designed as basements below ground surface. Therefore, more knowledge on how the ground shock propagates and attenuates during the scaled distances of 0.1 to 10 kg/m^{1/3} and effects on buried shelters, is needed.

During the Second World War, extensive experiment series and research were conducted on ground shock generated by high explosives [4]. This early work functions as a foundation for understanding the behaviour of how the shock waves propagates and attenuates in earth media. In [4], it is stated that the pressure in the soil from the detonation of an explosive charge is propagated by a plastic wave which is characterized by a continuous change of shape and of duration with distance from the charge, see Fig. 1 for schematic illustration. Close to the source the shape will be sharp and after a distance it will start to attenuate, and the wave will change shape.



Fig. 1 Schematic illustration of how the free field ground shock wave change shape with distance during propagation in a compactable soil.

When it comes to the structural response from the propagated pressure wave from ground shock it is not as obvious as when it comes to the airblast case. When predicting the structural response from airblast it is usually sufficient to separate the simulation in two parts:

- 1) Calculation of the pressures acting on rigid body, shaped as the studied structure, from e.g. an airblast simulation including the explosive and air formulation in a multi material Euler grid.
- 2) Determine the structural response by applying the calculated pressure time histories from 1) on the deformable structure modelled by e.g. shell elements.

This procedure is suitable for stiff structures, such as reinforced concrete structures, when subjected to air blast loading. For simple cases where the threat is directly in front of the structure step 1 above can be simplified to use empirical equations, e.g. for the incident airblast and adjust what the reflected pressure should be on the structure. However, this procedure is not applicable when it comes to the structural response due to ground shock simulations since, in ground shock, the reflected pressure actually depends on the movable mass and stiffness of the structure [5].

Hence, to determine the pressure acting on a structure, caused by ground shock, one must also know the properties of the structure. In 1979-1980 S. Hultgren at the National Fortifications Administration in Sweden, FortF, conducted experiments with a simplified structure, a buried suspended piston with mass m and stiffness k. Neglecting the influence of damping, the response of this structure can be described using a single degree of freedom system by using equation (1)

$$P_r(t) \cdot A = m \cdot a_x(t) + k \cdot d_x(t) \tag{1}$$

where $P_r(t)$ is the reflected pressure acting on the piston surface area A, a_x is the acceleration of the piston mass and d_x is the displacement of the piston.

These results show the same trend on how the reflecting pressure builds up by mass inertia and spring stiffness, see Fig. 2. Experimental results from reports [5]-[6] can be summed up in two points:

- If the structure has a high mass but low stiffness, then the reflected pressure time history results in a large first peak and a low second peak.
- If the structure has a low mass but high stiffness, the reflected pressure time history gets a low initial peak and a high second peak.

In 1985, experiments were conducted for buried deformable reinforced concrete walls with similar results [7]. This paper compare the experimental results from [5]-[6] with simulations carried out in AUTODYN-2D [8] and a simplified 2DOF model. Some comparisons have also been carried out in AUTODYN-2D [1]. These early simulations confirmed the main trends from the experiments. However, the symmetric 2D simulations in [1] included uncertainties such as improper air blast release due to buried explosive and uncertainty in actual reflected pressure measurements on the 2D piston. These uncertainties have been validated in 3D simulation model presented in paper [2]. The results in [2] conclude that the influence of not modelling surface release correctly in 3D is a marginal difference to the axi-symmetric 2D results for this experimental setup. This since the charge approximately can be regarded as fully buried. Therefore, the less computational expensive 2D axi-symmetry is used in this paper.



Fig. 2 Schematic illustration of how the reflection pressure $P_r(t)$ is made up of mass inertia $m \cdot a_x$ and stiffness $k \cdot d_x$ of a buried structure.

The paper is organized as follows: The section ORIGINAL HULTGREN EXPERIMENTAL SETUP AND EARLIER RESULTS, describes how the experiments were conducted and what parameters were varied. In section FE SIMULATION MODEL AND PARAMETERS STUDIED it is shown how the Euler and Lagrange elements were designed in AUTODYN-2D and what material models were used. In section 2DOF MODEL it is shown what parameters and initial conditions were used in the 2DOF model. In the section SIMULATION RESULTS the results from the extended parameter variation from obtained from simulations in AUTODYN-2D are analysed and compared with 2DOF. Finally, the section CONCLUSIONS AND FUTURE WORK concludes the findings from simulations and proposes suggestions for future work.

ORIGINAL HULTGREN EXPERIMENTAL SETUP AND EARLIER RESULTS

The experiments about the reflected pressure on a buried single degree of freedom system were conducted in sand. Both the charge of TNT with weight 0.5 kg and the suspended piston was buried 1 m, and the distance between the charge and the piston was set to 1 m. Unfortunately, the in-situ density of the sand was not measured nor the actual water content. However, in [5] it is mentioned that the sand was well compacted and not saturated. The experimental set up is shown in Fig. 3.



Fig. 3 Principal sketch of the experimental setup of [5]: buried piston with suspended mass m and stiffness k.

The main cylinder body consisted of a circular steel tube of 1.18 m in length and 0.36 m in outer diameter. The cylinder wall thickness was 8 mm and the diameter of the piston surface was 0.34 m. To increase the weight of the cylinder tube lead pieces were bolted to its inside; thus, increasing the total weight of the cylinder body to a total of 295 kg. The piston was movable on ball bearings through an axis and the stiffness was obtained with a helical spring.

The piston's movable mass was made of a removable plate and by changing plates with different thickness the suspended mass m of the piston was varied [5.2, 10.6, 24.7, and 58.8 kg]. Further, by changing the helical spring the stiffness k was also varied [0.1, 0.5, and 1.2 MN/m].

The following measurements were installed: piston accelerometer (a_x) , piston reflected pressure gauge (P_r) , and relative piston displacement (d_x) , see Fig. 3. Further details about the experimental setup, e.g. what model type of sensors was used, can be found in [5].

Hultgren's experimental results are presented in Fig. 4. When the piston mass was varied [5.2, 10.6, 24.7, and 58.8 kg], and the spring stiffness was kept constant to k = 0.1 MN/m, pressures according to Fig. 4 is shown in [5]. The dotted and dashed lines in Fig. 4 represent the estimated inertial pressure $P_{inertia}$ and spring pressure P_{spring} , defined as shown in equations (2) and (3), respectively. From this it can be noted that the dotted line in Fig. 4 follows well the measured pressure for the first peak; i.e. the first peak of the reflected pressure directly depends on the piston mass.





Fig. 4 Measured reflected pressure, estimated mass inertia pressure, and spring pressure for various piston masses and with constant spring stiffness k = 0.1 MN/m. From [5].

For the results in Fig. 5-Fig. 7 the spring stiffness was increased to k = 0.5 MN/m while the piston mass was varied [5.2, 10.6, 24.7, and 58.8 kg]. Now it can be noticed that the second ridge of the reflected pressure is more dominant in Fig. 5 compared to that in Fig. 4. A clear first peak is not visible until the piston mass reaches 24.7 kg. In Fig. 6 the reflected pressures from Fig. 5 are merged into one plot and from this it can be seen that when the piston mass was 5.2 and 10.6 kg there is no clear first peak in the reflected pressure. However, all of them have a second ridge in the reflected pressure, even though it is slightly reduced when the piston mass is increased.



Fig. 5 Measured reflected pressure, estimated mass inertia pressure, and spring pressure for various piston masses and with constant spring stiffness 0.5 MN/m. From [5].



Fig. 6 Comparison of measured reflected pressure from Fig. 5 for various piston masses and with constant spring stiffness 0.5 MN/m. From [5].

In Fig. 7a, the piston mass was kept constant at m = 5.2 kg while the spring stiffness was varied [0.1, 0.5, 1.2 MN/m]. It shows clearly how the second ridge increases with the spring stiffness. In comparison, the reflected pressure is shown in Fig. 7b when the piston mass is high, m = 58.8 kg, while the spring stiffness was varied [0.1, 0.5, 1.2 MN/m]. From this it can instead be noted that the pressure is only very little affected by the stiffness; i.e. the reflected pressure is mainly governed by the piston mass. Still the trend with the second ridge increases somewhat with increasing spring stiffness.



Fig. 7 *Measured reflected pressures for varied spring stiffness and with constant piston mass: (a) 5.2 kg, and (b) 58.8 kg. From [5].*

These main trends have already been confirmed with FE-simulations in [1], [2]. In the FE simulations in [2] the reflected pressure $P_r(t)$ of the piston was determined as

$$P_r(t) = \frac{m \cdot a_x(t) + k \cdot d_x(t)}{A} \tag{4}$$

and in Fig. 8 the variation of this pressure is shown in for a case with intermediate spring stiffness (k = 0.5 MN/m), and two different masses.



Fig. 8 Comparison of simulated reflected pressure for various piston masses and with constant spring stiffness 0.5 MN/m. Note that both diagrams show the same result but with different scaling on the vertical axis. For the corresponding experimental results, see Fig. 6. Results from [2].

From this it can be noted that a distinct initial pressure spike appears and that it is considerably larger when the mass is high (m = 58.8 kg) compared to when the mass is low (m = 5.2 kg). The appearance of such a distinct pressure peak is not in conjunction with the experiments; compare reflected pressure in Fig. 4 and Fig. 5. Further, the forming of a first (inertia) and second (stiffness) pressure peak, as schematically illustrated in Fig. 2, is not as clearly identified in the simulations as in the experiments. This is believed to, at least partly, be due to the strong initial pressure peak

obtained in the simulations. The reason for the above discrepancies are unknown. However, measuring the pressure in a material like sand is difficult and one possible reason might be that such an initial peak may have been filtrated away in the original experiments. The difference is noted but is apart from that not further treated here.

In Fig. 9 the reflected pressure from different simulations is compared for various stiffness when the piston mass is low (m = 5.2 kg) and high (m = 58.8 kg). When the piston mass is low it is clear that increased stiffness result in an increase of the second pressure peak (stiffness) in a way that is in good conceptual agreement with the experiments, compare with Fig. 7a. When the piston mass is high, though, the resulting reflected pressure is not as straightforward. Here, the simulations show that an increased stiffness still has a similar effect and thus cause an increased reflected pressure of the second pressure peak. However, this is not in full conjunction with the experimental observations, see Fig. 7b, where it is indicated that the effect of the stiffness is less important on the reflected pressure in the second pressure peak when the piston mass is high.



Fig. 9 *Simulated reflected pressures for varied spring stiffness and with constant piston mass 5.2 kg or 58.8 kg. For the corresponding experimental results, see Fig. 7. Results from [2].*



Fig. 10 Comparison of initial reflected pressure (solid lines) and reflected impulse intensity (dashed lines) obtained from 3D FE analyses and 2DOF model for various masses and spring stiffness 0.5 MN/m. Results from [2]. The resulting pressure from the 2DOF analysis showed that the first initial inertia pressure could be modelled. The pressure acting between sand cone and piston plate, is in Fig. 10 compared with the reflected pressure acting on the

piston plate in the 3D FE analyses for different masses when k = 0.5 MN/m. No influence on the resulting pressure was obtained due to spring stiffness; and hence it can be observed that for the cases studied the initial reflected pressure only depends on the piston mass. Further details about 3D FE-model setup or SDOF model setup can be found in [2]. Finally, if the incident wave was reflected against a rigid surface the pressure should increase with a factor of 2. Here, though, the ratio P_{load}/P_i is less than that; with decreasing value for decreasing piston mass. The reflection factor depends on the piston mass and varies from 0.97 to 1.98 (comparing with free field incident pressure in 2D) when the piston mass is increased from 5.2 to 58.8 kg, see [2].

FE SIMULATION MODEL AND PARAMETERS STUDIED

The FE model was built in AUTODYN [8] using a 2D axial symmetric geometry, see Fig. 11. The size of the model was 100 m x 50 m, which meant that the effects of reflections from the boundary could be entirely avoided for the piston response of at least 200 ms. There is a fine mesh rectangle 5.0 (6.0) m x 1.5 m containing the Free-Field target #25 and the right hand side of the fixed cylinder (orange) with square cell size of 2.5 mm. The sand, TNT, and VOID material inside the cylinder were modelled with multi-material Euler elements. The piston plate (blue) was modelled using rigid Lagrange elements constrained to move only in the piston direction. The cylinder body (orange) was also modelled using Lagrange elements and was fixed in all directions. The spring (purple) was modelled as a 1-element Lagrange spring with initial length of $L_0 = 10.0$ m. The purpose of the exaggerated spring length was to emulate an ideal linear Hook's law spring (F = k x) as close as possible ($L_0 >> x$). Using the 1-element Lagrange spring yields a spring stiffness of $k(x) = k_0 L_0/(L_0-x)$ with k_0 being the initial stiffness of the uncompressed spring.

$$F_{s}(t) = \frac{k_{0} \cdot L_{0}}{L_{0} - d(t)} d(t)$$
(5)

where *d* is the displacement of the plate (dark blue), see Fig. 11. The simulations were optimized for computational time by grid refinement using geometrical cell coarsening in the far field. The model was also parameterized so that charge size, distance to piston plate, piston plate mass, piston circular area, and spring stiffness could automatically be varied. Gravity was considered to have minimal impact on the simulation results and was hence not included in the model.



Fig. 11 2D axial symmetric model, with 25 fixed target points in the Eulerian SAND domain. The radial position of the four target rows i.e. 1-6, 7-12, 13-18, and 19-24 are located relative to the cylindrical plate (blue with radius r_{pl}). The 10m long Lagrangian spring (purple) is fixed at its right hand side and connected to the cylindrical plate. Target 26 is located on the piston plate.

Modelling of the sand

A simple way of modelling compaction materials is to define the plastic compaction curve as a pressure function of density $P(\rho)$ and the unloading wave speed as a function of density $c(\rho)$, see Fig. 12 for a schematic illustration. The unloading is then approximately represented by the straight blue lines on top of the dashed unloading curves. The solid straight line marked with TMD means Theoretical Maximum Density.

Earlier work on deriving mechanical properties for dry sand from tri-axial experiments, [10] and [11], does not include an easy change on how moisture affects the input parameters of the compaction EOS. When the soil type changes in water contents, porosity, and soil skeleton, the original compaction model [10] would need a new set of input. This has been studied in [16] where a total of 16 generic soil types was generated from dry sand to fully saturated clay. However, the main deficiency of the original model [15] for the compaction EOS is that the unloading phase is too simplified to properly model the shock wave propagation and the change in shape of the pressure wave as stated in [4]. In [18] an initial study was made of what the solution of unloading wave $c(\rho, P)$ for the dry sand would look like by using Non-Linear Programming (NLP) for numerically approximating each unloading curve. This allowed the experimental data from [10] and [11] to be extrapolated for the whole density pressure domain of interest [19].



Fig. 12 Schematic illustration of an EOS compaction, where the solid line between $(\rho_0, P_0 = 0)$ and (ρ_n, P_n) shows a plastic compaction curve, dashed red lines illustrate nonlinear unloading and straight blue lines show linear approximation with elastic unloading wave $c(\rho)$.

However, in this paper only the original (denoted Sjöbo) material model was used in the FE 2D-axisymmetry simulations, comparison of the modified material model (ModEOS) and Sjöbo was done in earlier 2D simulations [20]. In Sjöbo material model, the unloading speed only depends on density $c(\rho)$. The choice of not including both material models in the study was to focus on the main dynamic behaviour related to the interaction of the sand and the buried piston. Another benefit with Sjöbo sand is that original data and EOS and strength formulation is available for all Autodyn users in the standard library. The material properties for both volume and shear behaviour, and the user subroutines for the modified EOS and modified shear strength model for AUTODYN can be found in [20].

2DOF MODEL

The 2DOF model simplifies the experimental setup by the following assumptions, the sand in front of the structure is assumed to be modelled as an added mass, see Fig. 13 The explosive load is converted to either a force as function of time acting on the added mass or as a mass with an initial particle velocity on the added mass. In Fig. 14., the principal system of 2DOF is shown. The spring k_1 can only transfer compressive forces while spring k_2 can transfer forces in both directions. The initial velocity v_0 (i.e. the sand particle velocity U_{p0}) is used for describing the movement of the sand generated by the explosion.



Fig. 13 Illustration of how the 2DOF model is set up. The piston mass is connected to a second mass which is assumed to be an added mass from the sand and during compression the masses have a spring stiffness based on the bulk modulus of the sand.



Fig. 14 Illustration of 2DOF model used. Spring k_1 can only transfer compressive forces while spring k_2 can transfer forces in both directions.

The equations used to solve the 2DOF model can be written as

$$\begin{bmatrix} m_1 & 0\\ 0 & m_2 \end{bmatrix} \begin{bmatrix} a_{x1}\\ a_{x2} \end{bmatrix} + \begin{bmatrix} k_1 & -k_1\\ -k_1 & k_1 + k_2 \end{bmatrix} \begin{bmatrix} d_{x1}\\ d_{x2} \end{bmatrix} = \begin{bmatrix} F_1(t)\\ 0 \end{bmatrix}$$
(6)

where m_1 and m_2 is the sand mass and mass of the piston plate, respectively, k_1 is the stiffness of the sand cone and k_2 is the stiffness of the spring. Further \ddot{u}_1 and \ddot{u}_2 are accelerations and u_1 and u_2 are displacements of masses m_1 and m_2 , respectively. To correctly simulate that sand cannot transfer tensile forces the sand spring k_1 was modified so that only compressive forces could be transferred; spring k_2 though was linear elastic in both directions. The influence of damping was assumed to be negligible and is hence not included in the model. The added mass of sand is calculated as the mass of a cone between the charge point and the piston plate (i.e. the base of the cone) as

$$m_{sand} = A \cdot \frac{r}{3} \cdot \rho_{sand} \tag{7}$$

where A is the area of the piston plate, r is the charge distance and $\rho_{sand} = 1674 \text{ kg/m}^3$ is the sand in situ density.

The linear spring stiffness of the sand is approximated as

$$k_{sand} = \frac{K_{sand} \cdot A}{l_{cone}} \tag{8}$$

where K_{sand} is the sand bulk modulus and

$$l_{cone} = \frac{r}{3} \tag{9}$$

is the distance from the centre point of the sand cone to the piston plate. The bulk modulus was approximated to be equal both for loading and unloading cases and was determined as

$$K_{sand} = \rho_{sand} \cdot c_{sand}^2 \tag{10}$$

where $c_{sand} = 350$ m/s is the speed of the pressure wave (a typical value for dry soil [20]). This gives a sand bulk modulus of $K_{sand} = 0.2$ GPa (as a comparison, this is about ten times smaller than the bulk modulus of water, $K_{water} = 1000 \cdot 1484^2 = 2.2$ GPa). The model parameters for all cases are gathered in Table 1.

Table 1Model parameters used in the 2DOF model for the cases studied.

Identification		Part 1	: Sand	Part 2: Spring		
Group	Case	т	k	<i>m</i> sand	ksand	
		[kg]	[MN/m]	[kg]	[MN/m]	
1	1-12	5.2 - 58.8	0.1 - 1.2	50.7	56	
2	13-24	5.2 - 58.8	0.1 - 1.2	101.3	112	
3	25-36	5.2 - 58.8	0.1 - 1.2	202.6	223	
4	37-48	5.2 - 58.8	0.1 - 1.2	50.7	56	
5	49-60	5.2 - 58.8	0.1 - 1.2	50.7	56	
6	61-72	5.2 - 58.8	0.1 - 1.2	50.7	56	
7	73-84	5.2 - 58.8	0.1 - 1.2	50.7	56	
8	85-96	5.2 - 58.8	0.1 - 1.2	101.3	28	
9	97-108	5.2 - 58.8	0.1 - 1.2	202.6	223	

In this paper the hypothesis was that an effective sand mass, with an initial particle velocity, will give satisfying agreement with the FE results of initial reflected pressure for the different piston masses and spring deflection. In previous work, see [1] and [2], the initial particle velocity used was derived with ConWep [21] for a sand with 1674 kg/m³ in density, seismic wave speed 350 m/s, and 2.75 attenuation factor. This gave a particle velocity of 1.5 m/s for the selected charge size and distance to target. Here, though, particle velocities from Autodyn simulations have been used instead; for more detail see COMPARISON WITH 2DOF MODEL.

SIMULATION RESULTS

Simulations were carried out for a total of 12 combinations of different piston mass [5.2, 10.6, 24.7, 58.8 kg] and spring stiffness [0.1, 0.5, 1.2 MN/m]. For each of these combinations, the charge size W [0.5, 0.75, 1.0, 1.5 kg TNT] and the horizontal distance r, between charge and piston plate, was in one case also varied [1, 2 m]. Further, for some groups the piston contact surface was increased in size: $\Delta A = (A-A_0)/A_0 \cdot 100$, [0, 100, 300 %]. A total of 9 groups, each with 12 simulations, were conducted; resulting in a total number of 108 simulations, see Table 2 and APPENDIX I – CONDUCTED SIMULATION MATRIX.

Group	Case	W	r	ΔA
		[kg]	[m]	[%]
1	1-12	0.50	1.0	0
2	13-24	0.50	1.0	100
3	25-36	0.50	1.0	300
4	37-48	0.75	1.0	0
5	49-60	1.00	1.0	0
6	61-72	1.25	1.0	0
7	73-84	1.50	1.0	0
8	85-96	0.50	2.0	0
9	97-108	1.50	1.0	300

Table 2Summary of simulations carried out. In each group the piston mass [5.2, 10.6, 24.7, 58.8 kg] and spring
stiffness [0.1, 0.5, 1.2 MN/m] were varied, resulting in a total of 108 simulations.

The influence of the size of the piston surface is studied for the load case 0.5 kg TNT and r = 1 m during the initial phase up to 5 ms. In Fig. 15, the particle velocity in the sand just in front of the piston is studied in target points #1, #2, #3, see Fig. 11 for point positions. In addition, the piston velocity is also studied. All pistons have the same stiffness 0.1 MN/m, upper plots shows piston with mass 5.2 kg and lower plots with mass 58.8 kg. The increased piston surface increases from plot (a) 0 % to (b) 100 % and finally to (c) 300 %. For the piston with mass 5.2 kg there is only a marginal increase in piston velocity due to increase in piston area. However, For the piston with mass 58.8 kg both the response time is shortened from 1.8 ms to 0.5 ms and the piston speed is also increased from 1.1 m/s to 1.9 m/s during the initial 5 ms. This initial part of the piston movement is only mass inertia dependent, see Section ORIGINAL HULTGREN EXPERIMENTAL SETUP AND EARLIER RESULTS. When studying the ground shock pressure in front of the piston surface is increased, see upper plots in Fig. 16. When the piston with mass 58.8 kg is studied it can be seen that the amplitude of the pressure is the same, e.g. target point #1 reaches a pressure of 800 kPa, but the duration of the positive phase of ground shock decreases with increased piston surface, see upper plots in Fig. 16.



Fig. 15 Piston velocity #26 and particle velocity in sand #1, #2,#3 as a function of time. For load case 0.5 kg TNT and r=1.0 m. Upper plots (1), shows how piston weight 5.2 kg reacts for increased piston surface 0, 100, and 300%, (1a),(1b), and(1c), respectively, Lower plots (2), shows how piston weight 58.8 kg reacts for increased piston surface 0, 100, and 300%, (2a),(2b), and(2c), respectively. All cases have piston spring stiffness of 0.1 MN/m.



Fig. 16 Reflected pressure in sand #1, #2,#3 as a function of time. For load case 0.5 kg TNT and r = 1.0 m. Upper plots (1), shows how piston weight 5.2 kg reacts for increased piston surface 0, 100, and 300%, (1a),(1b), and(1c), respectively, Lower plots (2), shows how piston weight 58.8 kg reacts for increased piston surface 0, 100, and 300%, (2a),(2b), and(2c), respectively. All cases have piston spring stiffness of 0.1 MN/m.

In Fig. 17 and Fig. 18 the particle velocity and reflected pressure, respectively, are shown for the cases when the horizontal charge distance is kept constant but the charge weight is increased from 0.5 to 1.0 to 1.5 kg TNT. In the upper plots of Fig. 17 it is shown that the particle velocity, both before and after impact of piston with mass 0.5 kg and stiffness 0.1 MN/m, is steadily increasing from 2 to 3 to 4 m/s, respectively. In the lower plots of Fig. 17 it can be seen that the pressure, before impact of piston with mass 58.8 kg and stiffness 0.1 MN/m, is increasing as before. The velocity after impact, though, is lower: 1 to 2 to 2.5 m/s, respectively. In the upper plots of Fig. 18 one can see a steady increase in reflected pressure for piston with mass 5.2 kg. Here, the reflected maximum pressure increases from 600 to 900 to 1150 kPa. In the lower plots of Fig. 18 the results for the piston with mass 58.8 kg and stiffness 0.1 MN/m is shown. It can be seen that both the reflected maximum pressure and the positive phase duration is now higher compared to the piston with mass 5.2 kg. The maximum pressure increases from 800 to 1200 to 1500 kPa, and the positive duration increase from 0.3 ms (mass 5.2 kg) to 1.8 to 2.0.



Fig. 17 Piston velocity #26 and particle velocity in sand #1, #2,#3 as a function of time. For load cases 0.5, 1.0, and 1.5 kg TNT and r=1.0 m. Upper plots (1), shows how piston weight 5.2 kg reacts for increased charge weight 0.5, 1.0, and 1.5 kg TNT, (1a),(1b), and(1c), respectively, Lower plots (2), shows how piston weight 58.8 kg charge weight 0.5, 1.0, and 1.5 kg TNT, (2a),(2b), and(2c), respectively. All cases have piston spring stiffness of 0.1 MN/m.



Fig. 18 Reflected pressure in sand #1, #2,#3 as a function of time. For load cases 0.5,1.0, and 1.5 kg TNT and r=1.0 m. Upper plots (1), shows how piston weight 5.2 kg reacts for increased charge weight 0.5, 1.0, and 1.5 kg TNT, (1a),(1b), and(1c), respectively, Lower plots (2), shows how piston weight 58.8 kg charge weight 0.5, 1.0, and 1.5 kg TNT, (2a),(2b), and(2c), respectively. All cases have piston spring stiffness of 0.1 MN/m.

In Fig. 19 and Fig. 20 the particle velocity and reflected pressure, respectively, is shown for the cases when the horizontal charge distance is kept constant to 1.0 m and the charge weight is set to 1.5 kg TNT. The piston surface is increased from 0 to 300%, respectively. The upper plots in Fig. 19 show that for the piston with mass 5.2 kg the particle and piston velocity will be approximately the same, 4 m/s, with a slight increase for the after impact speed when the surface is 300%, see Fig. 19 (1b). Similar results for reflected pressure, compare (1a) and (1b) in Fig. 20, where pressure profiles are identical for 0 % and 300 % increased piston surface. When comparing the piston and the particle velocity for the piston with mass 58.8 kg and stiffness 0.1 MN/m. The piston and particle velocity after impact increases from 2.5 m/s to 3 m/s and above for the particle velocity in the soil when the surface is increased to 300%, compare (2a) and (2b) in Fig. 19. When the reflected pressure is compared for the piston with mass 58.8 kg and with stiffness 0.1 MN/m, it can be seen that increasing the surface reduces the positive phase duration of the pressure and slightly reduces the maximum reflected pressure of 1500 kPa, compare (2a) and (2b) in Fig. 20.



Fig. 19 Piston velocity #26 and particle velocity in sand #1, #2,#3 as a function of time. For load cases 1.5 kg TNT and r=1.0 m. Upper plots (1), shows how piston weight 5.2 kg reacts for increased piston surface 0 and 300 %, (1a) and (1b), respectively, Lower plots (2), shows how piston weight 58.8 kg reacts for increased piston surface 0 and 300 %, (1a) and (1b), respectively. All cases have piston spring stiffness of 0.1 MN/m.



Fig. 20 Reflected pressure in sand #1, #2,#3 as a function of time. For load cases 1.5 kg TNT and r=1.0 m. Upper plots (1), shows how piston weight 5.2 kg reacts for increased piston surface 0 and 300 %, (1a) and (1b), respectively, Lower plots (2), shows how piston weight 58.8 kg reacts for increased piston surface 0 and 300 %, (1a) and (1b), respectively. All cases have piston spring stiffness of 0.1 MN/m.

The total spring displacement is mainly influenced by the spring stiffness and then how big the piston area is, see Fig. 21. In addition the piston mass is only marginally influencing the final spring displacement. In Fig. 22 it is shown for the largest charge weight 1.5 kg TNT, that the piston area is giving several times larger final deflection when the surface area increased with 300%; e.g. the displacement was 107 mm for 5.2 kg piston mass and 0.1 MN/m with 0 % increased piston surface and 510 mm for 5.2 kg piston mass and 0.1 MN/m with 300 % increased piston surface. A steady increase for final displacement was seen when the charge weight was increased, see Fig. 23.

When it comes to the piston reflected pressure it is really built up from and calculated by using eq. (4), where the inertia pressure is mass dependent and the spring pressure part is stiffness dependent. Here, it is shown that when increasing the piston surface the initial inertia peak reflected pressure on the piston is reduced, see Fig. 24. When comparing the maximum deflection for the load cases with different charge weight, it can be seen that the displacement is steadily increased, see Fig. 25.Fig. 23



Fig. 21 Piston displacement #26 as a function of time. For load case 0.5 kg TNT and r=1.0 m. Upper plots (1), shows how piston weight 5.2 kg reacts for increased piston surface 0, 100, and 300%, (1a),(1b), and(1c), respectively, Lower plots (2), shows how piston weight 58.8 kg reacts for increased piston surface 0, 100, and 300%, (2a),(2b), and(2c), respectively. All plots show how piston spring stiffness of 0.1 0.5 and 1.2 MN/m influences the final displacement.



Fig. 22 Piston deflection #26 as a function of time. For load cases 1.5 kg TNT and r=1.0 m. Upper plots (1), shows how piston weight 5.2 kg reacts for increased piston surface 0 and 300 %, (1a) and (1b), respectively, Lower plots (2), shows how piston weight 58.8 kg reacts for increased piston surface 0 and 300 %, (1a) and (1b), respectively. All plots show how piston spring stiffness of 0.1 0.5 and 1.2 MN/m influences the final displacement.



Fig. 23 Piston deflection #26 as a function of time. For load cases 0.5,1.0, and 1.5 kg TNT and r=1.0 m. Upper plots (1), shows how piston weight 5.2 kg reacts for increased charge weight 0.5, 1.0, and 1.5 kg TNT, (1a),(1b), and(1c), respectively, Lower plots (2), shows how piston weight 58.8 kg charge weight 0.5, 1.0, and 1.5 kg TNT, (2a),(2b), and(2c), respectively. All plots show how piston spring stiffness of 0.1 0.5 and 1.2 MN/m influences the final displacement.



Fig. 24 Piston reflected pressure #26 as a function of time. For load case 0.5 kg TNT and r=1.0 m. Upper plots (1), shows how piston weight 5.2 kg reacts for increased piston surface 0, 100, and 300%, (1a),(1b), and(1c), respectively, Lower plots (2), shows how piston weight 58.8 kg reacts for increased piston surface 0, 100, and 300%, (2a),(2b), and(2c), respectively.



Fig. 25 Piston reflected pressure #26 as a function of time. For load cases 0.5,1.0, and 1.5 kg TNT and r=1.0 m. Upper plots (1), shows how piston weight 5.2 kg reacts for increased charge weight 0.5, 1.0, and 1.5 kg TNT, (1a),(1b), and(1c), respectively, Lower plots (2), shows how piston weight 58.8 kg charge weight 0.5, 1.0, and 1.5 kg TNT, (2a),(2b), and(2c), respectively.

Fig. 26 shows pressure plots from the ground shock arriving at the piston and it can be seen that the initial reflection part is already over at 3.5 ms. This has also been confirmed in earlier studies [2].



Fig. 26 Pressure plots at different times from 2D-FE simulations of 58.8 kg and 0.1 MN/m, with charge weight 1.0 kg with horizontal distance 1 m, Blue = 0 kPa, green = 750 kPa, red \geq 1500 kPa.

Fig. 27 shows the shear stress (ZX) for the same times as in Fig. 26 and indicates that in front of the piston there is a cone shaped part which can be seen at 3 ms. This can be an indication that there is an effective sand mass, i.e. the sand inside the cone, colliding with the piston at a certain particle velocity.



Fig. 27 Shear stress (ZX) plots at different times from times from 2D-FE simulations of 58.8 kg and 0.1 MN/m, with charge weight 1.0kg with horizontal distance 1 m, Blue \leq -200 kPa, green = 0 kPa, red \geq 200 kPa.

In Fig. 28. the corresponding sand particle velocity $U_{p,x}$ is shown. Here it can be noted that the sand velocity directly in front of the piston is larger compared to the sand outside the piston; i.e. these plots indicate that the sand in front of the projected area of the piston encounter less resistance compared to the sand outside this zone. This response is conceptually very different to that obtained in e.g. a shock wave in air hitting an object, in which the shock wave in front of the object instead would reflect against it. The shape of the velocity distribution differs from the cone shaped part found in Fig. 27. This indicates that the volume of the effective mass colliding with the piston may be larger than that assumed in the 2DOF simulations. It seems that the mass transportation away from the centre of the explosive detonation increases in the direction of the piston, and the simple reason is that the spring stiffness for the piston is maximum 1.2 kN/m; i.e. more than 45 times smaller than the sand stiffness of 56 kN/m, see Table 1. This gives a reason to transport more sand material towards the piston area, which is more easily compressed than the surrounding sand. This is also confirmed in the t = 10 ms plot in Fig. 28.



Fig. 28 Absolute velocity plots at different times from plots at different times from times from 2D-FE simulations of 58.8 kg and 0.1 MN/m, with charge weight 1.0kg with horizontal distance 1 m, Blue ≤ 0 m/s, green = 1.5 m/s, red ≥ 3 m/s.

COMPARISON WITH 2DOF MODEL

Comparisons of Autodyn results have been made with a 2DOF model. Here, the initial velocity of the piston plate, v_{plate} , and the maximum internal energy of the spring, W_i , is compared. The initial plate velocity is based on classic elastic impact theory and can be determined as

$$v_{plate} = \frac{2m_{sand}}{m_{sand} + m_{plate}} \cdot v_{imp} \tag{11}$$

where $v_{imp} = U_{p,imp}$ is the particle velocity of the sand at first impact. In the 2DOF analyses it was noticed in [2] that the spring internal energy converges towards the kinetic energy of the incoming mass; i.e.

$$W_{i,2DOF} = E_{k,sand} = \frac{m_{sand} \cdot v_0^2}{2}$$
(12)

where the velocity $v_0 = U_{p,sand}$ is a velocity that represent a energy weighted mean velocity of the whole sand cone. Hence, $U_{p,sand} \neq U_{p,imp}$; instead it is a velocity that relates to the kinetic total energy of the sand cone just prior to impact to the piston plate. This weighted velocity can be determined as

$$U_{p,sand} = \frac{\sum \left(\Delta m_{sand,i} \cdot U_{p,i}^{2} \right)}{m_{sand}}$$
(13)

where $\Delta m_{sand,i}$ and $U_{p,i}$ is the mass and particle velocity, respectively, in part *i* of the sand cone at the time of impact.

The spring internal energy can be determined as

$$W_i = \frac{k \cdot d_{x^2}}{2} \tag{14}$$

where k is the spring stiffness and d_{x2} is the deformation of the spring.

In this comparison particle velocities from the Autodyn simulations have been used rather than particle velocities determined using ConWep [21]; see Fig. 29 for an example of $U_p(t)$ at different distances from the charge in a free field environment. The sand particle velocities at impact in ConWep and Autodyn are compared in Fig. 30. From this it can be noted that the difference is rather substantial in most cases. This may be due to the choice of attenuation factor used in ConWep (2.75 was used) but is not further dealt with here. Further, the weighted particle velocity $U_{p,sand}$, according to equation (13), is also included in order to determine the total kinetic energy of the sand cone.



Fig. 29 Particle velocity at different distances from the charge for simulations in Group 1 (W = 0.5kg TNT, r = 1.0 m).



Fig. 30 Comparison of particle velocities in ground obtained using ConWep and the present Autodyn analyses. The effective particle velocity $U_{p,sand}$, describes an average velocity in the sand cone used to determine its kinetic energy at the time the shock front reaches the piston plate, see equation (13).

In Fig. 31 a comparison of the velocity ratio

$$\eta_{v} = \frac{v_{plate,Autodyn}}{v_{plate,2DOF}}$$
(15)

and energy ratio

$$\eta_W = \frac{W_{i,Autodyn}}{W_{i,2DOF}} \tag{16}$$

is shown for all 108 simulations. From this it can be seen that the initial plate velocities is rather similar in Autodyn and 2DOF; the ratios mostly being within 0.8-1.1. For the internal energy, though, the conformance is not as good, even though the ratios for most cases are within 0.6-1.5. It can be noted that a low stiffness and large piston plate area generates a larger value on η_W while, the plate mass, though, seems to have a rather minor effect on the resulting ratio (only load cases 85-96, i.e. Group 8, is affected much). Further, an increased charge weight somewhat decreases the energy ratio η_W . For load cases with the charge located at distance r = 2.0 m from the piston plate, i.e. Group 8, the largest deviations in both velocity and energy ratios are observed. The reason for this is unknown and need further studies.



Fig. 31 Comparison of velocity and energy ratios of the plate and spring stiffness obtained using Autodyn and 2DOF for all 108 cases studied.

CONCLUSIONS AND FUTURE WORK

2D Finite Element simulations have been carried out in Autodyn to study the structural response of a well-defined structure; a suspended piston-spring system buried in sand subjected to ground shock from an explosive charge. The parameters varied in the simulations were charge size, charge distance, reflection area of the piston, piston mass, and spring stiffness; a total of 108 load cases were studied. Further, an attempt has been made to approximately describe the structural response using a simplified 2DOF model; describing the sand and structure, respectively.

A simplified calculation method, based on a 2DOF model, has been used to approximately describe the initial and final structural response of the loaded structure. The initial response is described using classic elastic impact theory, based on sand particle velocity in shock front, and the final response is based on energy balance of the total external kinetic energy of the sand cone at time of impact and the internal energy consumption of the linear elastic spring. The comparisons made show that this simplified approach rather well describes the results obtained in the FE simulations of the cases studied here. However, there are still need of further improvement and comparisons of this approach. Based on the cases studied in here it can be noted that the prediction ability of the simplified method decreases when the load area increases and that the simplified model is not able to adequately consider the influence of spring stiffness. Further, it is implied (only one group studied) that an increased distance to the charge decrease the success level to accurately predict the final structural response.

Further studies will be made to examine if the simplified approach can be further improved. As part of this the number of load cases studied using FE simulations will also increase; with a special focus on the response obtained at increased charge distances.

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REFERENCES

- [1] Laine, L., Johansson, M., and Larsen O.P. (2015): "Simulation of experiments which show that reflection pressure time history from ground shock depends on the reflected structure's stiffness and mass", *Proceedings of the 86th Shock and Vibration Symposium*, Orlando, FL, USA.
- [2] Laine, L., Johansson, M., and Larsen O.P. (2016): "3D FE and 2DOF simulations of ground shock experiments – Reflection pressure time history dependency due to the structure's stiffness and mass", *Proceedings of the 87th Shock and Vibration Symposium*, New Orleans, LA, USA.
- [3] Ekengren B. (2015): *Skyddsrum, SR 15* (Civil Defence Shelters SR 15, in Swedish.), the Swedish Civil Contingencies Agency (MSB), report no. MSB748, ISBN 978-91-7383-485-8, Karlstad, Sweden.
- [4] Lampson C.W. (1946): *Final Report on Effects of Underground Explosions*, Div. 2, National Defence Research Committee of the US Office Scientific R&D, NDRC Report No. A-479, OSRD Report No. 6645.
- [5] Hultgren S. (1979): *Explosion of buried model structures to buried TNT explosions in sand*, Fortfikationsförvaltningen Fortf, Report nr C 183, Eskilstuna, Sweden.
- [6] Hultgren S. (1980): (*Effect of underground explosions in the sand with deformable walls, in Swedish*), Omgång II, Fortfikationsförvaltningen Fortf, Rapport nr C 200, Eskilstuna, Sweden.
- [7] Hultgren S. (1985): On the effects of buried explosions in sand, Fortfikationsförvaltningen Fortf, Report nr C2:85, Eskilstuna, Sweden.
- [8] Century Dynamics Inc. (2004): AUTODYN Theory Manual Revision 5.0, San Ramon, CA, USA.
- [9] Johansson, M., Larsen O.P., and Laine, L. (2007): "Explosion at an intersection in an Urban Environment Experiments and analyses", *Proceedings of the 78th Shock and Vibration Symposium*, Philadelphia, PA, USA.
- [10] Laine L. and Sandvik A. (2001): "Derivation of mechanical properties for sand", 4th Asian-Pacific conference on Shock and Impact Loads on Structures, CI-Premier PTE LTD, vol. 4, pp 353-360, Singapore.
- [11] Heyerdahl H. and Madshus C. (2000): "EOS-data for sand, Tri-axial tests on sand from Sjöbo", Norges Geotekniske institutt, NGI rept. 20001157-1, Oslo, Norway.
- [12] Fairlie G. and Bergeron D. (2002): "Numerical simulations of Mine Blast Loading on Structures", 17th Numerical aspects of Blast Symposium, Las Vegas, Nevada.
- [13] Tjernberg A. (2006): Simulation of Mine-Blast deflection, FOI-Swedish Defence Research Agency, Technical Report, FOI-R--1913--SE, TUMBA, Sweden.
- [14] Grujicic M., Pandurangan B., Qiao R., Cheeseman B.A., Roy W.N., Skaggs R.R., and Gupta R. (2008): "Parameterization of the porous-material model for sand with different levels of water saturation", *Soil Dynamics and Earthquake Engineering* 28, pp 20–35.
- [15] Moxnes J. F., Ødegårdstuen G., Atwood A., and Curran P. (1999): "Mechanical properties of a porous material studied in a high speed piston driven compaction experiment", 30th International Annual Conference of ICT Energetic Materials, Fraunhofer Institut Chemische Technologie.
- [16] Wang1 Z., Hao H., and Lu Y. (2004): "A three-phase soil model for simulating stress wave propagation due to blast loading", *Int. Journal for Numerical and Analytical Methods in Geomechanics*, 28:33–56 (DOI: 10.1002/nag.325).
- [17] Laine L. (2006): "Study of Planar Ground Shock in Different Soils and Its Propagation Around a Rigid Block", 77th Shock and Vibration Symposium, Monterey, CA.
- [18] Laine L. and Larsen O.P. (2009): "Proposal on How to Model the Unloading in a Compaction Equation of State based upon Tri-axial tests on Dry Sand", *80th Shock & Vibration Symposium*, San Diego, CA.
- [19] Laine L. and Larsen O.P. (2012): "Implementation of Equation of State for Dry Sand in Autodyn", 83rd Shock and Vibration Symposium, Shock and Vibration Exchange (SAVE), New Orleans, LA.
- [20] Laine L. (2012): Markstötvåg (Ground Shock. In Swedish). MSB, Myndigheten för samhällsskydd och beredskap. Publ.nr MSB344, Karlstad. Weblink:https://www.msb.se/Upload/Insats_och_beredskap/Olycka_kris/Skyddsrum/Litteratur/L01.%20Refer enslitteratur/L01-202 Markst%C3%B6tv%C3%A5g.pdf
- [21] ConWep (1992): Collection of conventional weapons effects calculations based on TM 5-855-1, Fundamentals of Protective Design for Conventional Weapons, U.S. Army Engineer Waterways Experiment Station, Vicksburg, USA.

Buried Plate-Spring System 2D Autodyn Simulations 2017								
	Run Status			Structure		Charge and Distance		Plate Area Increase
Sim ID	LowRez	HighRez	Group	m [kg]	k ₀ [M N/m]	W [kg TNT]	r [m]	(A-A0)/A0*100 [%]
c001	Х	Х	1	5.2	0.1	0.5	1	0
c002	Х	Х			0.5			
c003	Х	Х			1.2			
c004	Х	Х		10.6	0.1			
c005	Х	Х			0.5			
c006	Х	Х			1.2			
c007	Х	х		24.7	0.1			
c008	Х	Х			0.5			
c009	Х	Х			1.2			
c010	X	Х		58.8	0.1			
c011	X	X			0.5			
c012	Х	Х			1.2			
c013		х	2	5.2	0.1	0.5	1	100
c014		Х			0.5			
c015		Х			1.2			
c016		Х		10.6	0.1			
c017		Х			0.5			
c018		х			1.2			
c019		х		24.7	0.1			
c020		х			0.5			
c021		х			1.2			
c022		Х		58.8	0.1			
c023		Х			0.5			
c024		х			1.2			
c025		Х	3	5.2	0.1	0.5	1	300
c026		Х			0.5			
c027		Х			1.2			
c028		Х		10.6	0.1			
c029		Х			0.5			
c030		Х			1.2			
c031		х		24.7	0.1			
c032		Х			0.5			
c033		Х			1.2			
c034		Х	1	58.8	0.1			
c035		Х			0.5			
c036		Х			1.2			

APPENDIX I – CONDUCTED SIMULATION MATRIX

Buried Plate-Spring System 2D Autodyn Simulations 2017								
	Run Status			Structure		Charge and Distance		Plate Area Increase
Sim	LowRez	HighRez	Group	m	k ₀	W	r	(A-A0)/A0*100
ID				[kg]	[M N/m]	[kg TNT]	[m]	[%]
c037		Х	4	5.2	0.1	0.75	1	0
c038		Х			0.5			
c039		Х			1.2			
c040		Х		10.6	0.1			
c041		Х			0.5			
c042		Х			1.2			
c043		Х		24.7	0.1			
c044		Х			0.5			
c045		Х			1.2			
c046		Х		58.8	0.1			
c047		X			0.5			
c048		Х			1.2			
c049		Х	5	5.2	0.1	1	1	0
c050		Х			0.5			
c051		Х			1.2			
c052		Х		10.6	0.1			
c053		Х			0.5			
c054		Х			1.2			
c055		Х		24.7	0.1			
c056		Х			0.5			
c057		Х			1.2			
c058		Х		58.8	0.1			
c059		Х			0.5			
c060		Х			1.2			
c061	1	х	6	5.2	0.1	1.25	1	0
c062	1	х	1		0.5			
c063	1	X	1		1.2			
c064	1	х	1	10.6	0.1			
c065		X			0.5			
c066	1	X	1		1.2			
c067		X		24.7	0.1			
c068		x			0.5			
c069	1	X	1		1.2			
c070		X		58.8	0.1			
c071		X			0.5			
c072		X	1		1.2			

Buried Plate-Spring System 2D Autodyn Simulations 2017								
	Run Status			Structure		Charge and Distance		Plate Area Increase
Sim	LowRez	HighRez	Group	m	k ₀	W	r	(A-A0)/A0*100
ID				[kg]	[M N/m]	[kg TNT]	[m]	[%]
c073		Х	7	5.2	0.1	1.5	1	0
c074		Х			0.5			
c075		Х			1.2			
c076		Х		10.6	0.1			
c077		Х			0.5			
c078		Х			1.2			
c079		Х		24.7	0.1			
c080		Х			0.5			
c081		Х			1.2			
c082		Х		58.8	0.1			
c083		X			0.5			
c084		Х			1.2			
c085	X	Х	8	5.2	0.1	0.5	2	0
c086		Х			0.5			
c087		Х			1.2			
c088		Х		10.6	0.1			
c089		Х			0.5			
c090		Х			1.2			
c091		Х	-	24.7	0.1			
c092		Х	-		0.5			
c093		Х			1.2			
c094		Х		58.8	0.1			
c095		Х	-		0.5			
c096		х	-		1.2			
c097	X	X	9	5.2	0.1	1.5	1	300
c098		Х	-		0.5			
c099	1	X	1		1.2			
c100	1	X	1	10.6	0.1			
c101	1	X	1		0.5			
c102	1	X	1		1.2			
c103		X	-	24.7	0.1			
c104		x			0.5			
c105	1	X	1		1.2			
c106	X	X	-	58.8	0.1			
c107		x	-		0.5			
c108	1	x	1		1.2			